

# Nonlinear Formulation of DFT

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## 1 Introduction

In this text, we explore the mathematical structure of the DFT self-consistency cycle and show that it is nothing else than a standard nonlinear problem, only the Jacobian matrix is dense, so we need to use nonlinear solvers that can handle that.

## 2 Problem

The task is to find such a charge density  $n$ , so that all the equations below hold (e.g. are self-consistent):

$$V = -\frac{Z}{r} + V_H + V_{xc}$$

$$\left(-\nabla^2 + V\right) \phi_m = \epsilon_m \phi_m, \quad m = 1, 2, \dots, 4$$

$$n = \sum_{m=1}^4 \phi_m^2$$

$$V_{xc} = f(n)$$

$$\nabla^2 V_H = n$$

## 3 Reformulation

Let's write everything in terms of  $\phi_m(x)$  explicitly:

$$n(x) = \sum_{m=1}^4 \phi_m^2(x)$$

$$\begin{aligned}
V_{xc}(x) &= f(n(x)) = f\left(\sum_{m=1}^4 \phi_m^2(x)\right) \\
V_H(x) &= \frac{1}{2} \int_{\Omega} \frac{n(x')}{|x' - x|} \mathbb{X}' = \frac{1}{2} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} \mathbb{X}' \\
V(x) &= -\frac{Z}{r} + V_H(x) + V_{xc}(x) = \\
&= -\frac{Z}{r} + \frac{1}{2} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} \mathbb{X}' + f\left(\sum_{m=1}^4 \phi_m^2(x)\right)
\end{aligned}$$

Now we can write everything as just one (nonlinear) equation:

$$\left(-\nabla^2 - \frac{Z}{r} + \frac{1}{2} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} \mathbb{X}' + f\left(\sum_{m=1}^4 \phi_m^2(x)\right)\right) \phi_n = \epsilon_n \phi_n, \quad n = 1, 2, \dots, 4$$

## 4 FE Discretization

The correspondig discrete problem has the form

$$\begin{aligned}
\int_{\Omega} \nabla \phi_n(x) \cdot \nabla v_i(x) + \left[-\frac{Z}{r} + \frac{1}{2} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} \mathbb{X}' + f\left(\sum_{m=1}^4 \phi_m^2(x)\right)\right] \phi_n(x) v_i(x) \mathbb{X} &= \\
= \int_{\Omega} \epsilon_n \phi_n(x) v_i(x) \mathbb{X}, \quad n = 1, 2, \dots, 4; \quad i = 1, 2, \dots, N
\end{aligned}$$

where

$$\phi_n = \phi_n(\mathbf{Y}^{(n)}) = \sum_{j=1}^N y_j^{(n)} v_j(x)$$

Here  $\mathbf{Y}^{(n)} = (y_1^{(n)}, y_2^{(n)}, \dots, y_N^{(n)})^T$  is the vector of unknown coefficients for the  $n$ th wavefunction  $\phi_n(x)$ . Our equation can then be written in the compact form

$$\mathbf{F}(\mathbf{Y}^{(n)}) = \mathbf{0}, \quad n = 1, 2, \dots, 4$$

where  $\mathbf{F} = (F_1, F_2, \dots, F_N)^T$  with

$$\begin{aligned}
F_i(\mathbf{Y}^{(n)}) &= \int_{\Omega} \nabla \phi_n(x) \cdot \nabla v_i(x) + \left[-\frac{Z}{r} + \frac{1}{2} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} \mathbb{X}' + f\left(\sum_{m=1}^4 \phi_m^2(x)\right)\right] \phi_n(x) v_i(x) \mathbb{X} - \\
&\quad - \int_{\Omega} \epsilon_n \phi_n(x) v_i(x) \mathbb{X}
\end{aligned}$$

## 5 Jacobian

The Jacobi matrix has the elements:

$$J_{ik} = \frac{\partial F_i}{\partial y_k^{(s)}}$$

The only possible dense term is:

$$\begin{aligned} & \frac{\partial}{\partial y_k^{(s)}} \int_{\Omega} \int_{\Omega} \frac{\sum_{m=1}^4 \phi_m^2(x')}{|x' - x|} v_k(x') v_i(x) dx' dx = \\ &= \frac{\partial}{\partial y_k^{(s)}} \int_{\Omega} \int_{\Omega} \frac{\sum_{m=1}^4 \left( \sum_{j=1}^N y_j^{(m)} v_j(x') \right)^2}{|x' - x|} dx' \left( \sum_{j=1}^N y_j^{(n)} v_j(x) \right) v_i(x) dx = \\ &= \int_{\Omega} \int_{\Omega} \frac{2 \left( \sum_{j=1}^N y_j^{(s)} v_j(x') \right) v_k(x')}{|x' - x|} dx' \left( \sum_{j=1}^N y_j^{(n)} v_j(x) \right) v_i(x) dx + \\ &+ \int_{\Omega} \int_{\Omega} \frac{\sum_{m=1}^4 \left( \sum_{j=1}^N y_j^{(m)} v_j(x') \right)^2}{|x' - x|} dx' \delta_{ns} v_k(x) v_i(x) dx \end{aligned}$$

Now we can see that we have in there the following term:

$$\int_{\Omega} \int_{\Omega} \frac{v_k(x') v_i(x)}{|x' - x|} dx' dx$$

which is dense in  $(ki)$ , as can be easily seen by fixing  $i$  and writing

$$\int_{\Omega} \int_{\Omega} \frac{v_k(x')}{|x' - x|} dx' v_i(x) dx$$

so for each  $k$  there is some contribution from the integral  $\int_{\Omega} \frac{v_k(x')}{|x' - x|} dx'$  for such  $x$  where  $v_i(x)$  is nonzero, thus making the Jacobian  $J_{ik}$  dense.