

Taylor-Galerkin Method for Compressible Euler Equations

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January 2009

1 Euler Equations

$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \mathbf{v}) = 0 \quad (1)$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) + \nabla p = 0 \quad (2)$$

$$\frac{\partial \rho e}{\partial t} + \operatorname{div}((\rho e + p) \mathbf{v}) = 0 \quad (3)$$

Using the perfect gas state equation $p = \varrho RT$ with $R = c_p - c_v$, we obtain

$$p = (\kappa - 1) \left(\varrho e - \frac{1}{2} \varrho |\mathbf{v}|^2 \right),$$

where $\kappa = c_p/c_v$.

2 Equations in conservative variables

The conservative variables are defined as follows,

$$\mathbf{u} = \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{pmatrix} = \begin{pmatrix} \varrho \\ \varrho v_1 \\ \varrho v_2 \\ \varrho e \end{pmatrix},$$

where ϱ, v_1, v_2, e are the density, first velocity component, second velocity component, and total energy, respectively. In the conservation form, Euler equations can be written as

$$\frac{\partial \mathbf{u}}{\partial t} + \frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}) = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad (4)$$

where Eulerian fluxes have the following form

$$f_1(\mathbf{u}) = \begin{pmatrix} u_2 \\ \frac{u_2^2}{u_1} + (\kappa - 1)(u_4 - \frac{u_2^2 + u_3^2}{2u_1}) \\ \frac{u_3 u_2}{u_1} \\ \frac{u_2}{u_1} \left(\kappa u_4 - (\kappa - 1) \frac{u_2^2 + u_3^2}{2u_1} \right) \end{pmatrix}, \quad f_2(\mathbf{u}) = \begin{pmatrix} u_3 \\ \frac{u_2 u_3}{u_1} \\ \frac{u_3^2}{u_1} + (\kappa - 1)(u_4 - \frac{u_2^2 + u_3^2}{2u_1}) \\ \frac{u_3}{u_1} \left(\kappa u_4 - (\kappa - 1) \frac{u_2^2 + u_3^2}{2u_1} \right) \end{pmatrix},$$

Equation (4) can be written as

$$\frac{\partial \mathbf{u}}{\partial t} + A_1(\mathbf{u}) \frac{\partial \mathbf{u}}{\partial x_1} + A_2(\mathbf{u}) \frac{\partial \mathbf{u}}{\partial x_2} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad (5)$$

where the Jacobi matrices of the Eulerian fluxes \mathbf{f}_1 and \mathbf{f}_2 are defined as

$$A_1(\mathbf{u}) = D\mathbf{f}_1/D\mathbf{u}, \quad A_2(\mathbf{u}) = D\mathbf{f}_2/D\mathbf{u},$$

where Jacobi matrices $A_1(\mathbf{u})$ and $A_2(\mathbf{u})$ have the form

$$A_1(\mathbf{u}) = \begin{pmatrix} 0 & 1 & 0 & 0 \\ \frac{(\kappa - 1)(u_2^2 + u_3^2)}{2u_1^2} - \frac{u_2^2}{u_1^2} & (3 - \kappa) \frac{u_2}{u_1} & (1 - \kappa) \frac{u_3}{u_1} & \kappa - 1 \\ -\frac{u_2 u_3}{u_1^2} & \frac{u_3}{u_1} & \frac{u_2}{u_1} & 0 \\ \frac{u_2}{u_1} \left(\frac{(\kappa - 1)(u_2^2 + u_3^2)}{u_1^2} - \kappa \frac{u_4}{u_1} \right) & \kappa \frac{u_4}{u_1} - \frac{(\kappa - 1)u_2^2}{u_1^2} - \frac{(\kappa - 1)(u_2^2 + u_3^2)}{2u_1^2} & \frac{(1 - \kappa)u_2 u_3}{u_1^2} & \kappa \frac{u_2}{u_1} \end{pmatrix}$$

$$A_2(\mathbf{u}) = \begin{pmatrix} 0 & 0 & 1 & 0 \\ -\frac{u_2 u_3}{u_1^2} & \frac{u_3}{u_1} & \frac{u_2}{u_1} & 0 \\ \frac{(\kappa - 1)(u_2^2 + u_3^2)}{2u_1^2} - \frac{u_3^2}{u_1^2} & (1 - \kappa) \frac{u_2}{u_1} & (3 - \kappa) \frac{u_3}{u_1} & \kappa - 1 \\ \frac{u_3}{u_1} \left(\frac{(\kappa - 1)(u_2^2 + u_3^2)}{u_1^2} - \kappa \frac{u_4}{u_1} \right) & \frac{(1 - \kappa)u_2 u_3}{u_1^2} & \kappa \frac{u_4}{u_1} - \frac{(\kappa - 1)u_3^2}{u_1^2} - \frac{(\kappa - 1)(u_2^2 + u_3^2)}{2u_1^2} & \kappa \frac{u_3}{u_1} \end{pmatrix}$$

3 Taylor-Galerkin Method

Let us expand \mathbf{u} in time using second-order Taylor series

$$\mathbf{u}^{n+1} = \mathbf{u}^n + \Delta t \frac{\partial \mathbf{u}^n}{\partial t} + \frac{\Delta t^2}{2} \frac{\partial}{\partial t} \frac{\partial \mathbf{u}^{n+1}}{\partial t} + O(\Delta t^3). \quad (6)$$

Now, plug in the equation (4) to the Taylor series (6)

$$\frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\Delta t} + \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}^n) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}^n) \right) + \frac{\Delta t}{2} \frac{\partial}{\partial t} \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}^{n+1}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}^{n+1}) \right) = 0 \quad (7)$$

We can rewrite the third term in (7) in terms of spatial derivatives using integration by parts

$$\begin{aligned} \frac{\partial}{\partial t} \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}) \right) &= \frac{\partial}{\partial x_1} \frac{\partial}{\partial t} \mathbf{f}_1(\mathbf{u}) + \frac{\partial}{\partial x_2} \frac{\partial}{\partial t} \mathbf{f}_2(\mathbf{u}) = \frac{\partial}{\partial x_1} \frac{\partial \mathbf{f}_1}{\partial \mathbf{u}} \frac{\partial \mathbf{u}}{\partial t} + \frac{\partial}{\partial x_2} \frac{\partial \mathbf{f}_2}{\partial \mathbf{u}} \frac{\partial \mathbf{u}}{\partial t} = \\ &= -\frac{\partial}{\partial x_1} A_1(\mathbf{u}) \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}) \right) - \frac{\partial}{\partial x_2} A_2(\mathbf{u}) \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}) \right) \end{aligned} \quad (8)$$

Now, substituting (8) into (7) we obtain

$$\frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\Delta t} + \sum \frac{\partial}{\partial x_k} \mathbf{f}_k(\mathbf{u}^n) - \frac{\Delta t}{2} \sum \frac{\partial}{\partial x_k} A_k(\mathbf{u}^{n+1}) \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}^{n+1}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}^{n+1}) \right) = 0 \quad (9)$$

Proceeding with the Galerkin method, we multiply (9) by a test function \mathbf{v} and integrate over the domain Ω to get the following scheme

$$\begin{aligned} \int_{\Omega} \frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\Delta t} \mathbf{v} - \int_{\Omega} \sum \mathbf{f}_k(\mathbf{u}^n) \frac{\partial \mathbf{v}}{\partial x_k} + \int_{\partial \Omega} \sum \mathbf{f}_k(\mathbf{u}^n) \mathbf{v} n_k + \\ + \frac{\Delta t}{2} \int_{\Omega} \sum A_k(\mathbf{u}^{n+1}) \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}^{n+1}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}^{n+1}) \right) \frac{\partial \mathbf{v}}{\partial x_k} - \\ - \frac{\Delta t}{2} \int_{\partial \Omega} \sum A_k(\mathbf{u}^{n+1}) \left(\frac{\partial}{\partial x_1} \mathbf{f}_1(\mathbf{u}^{n+1}) + \frac{\partial}{\partial x_2} \mathbf{f}_2(\mathbf{u}^{n+1}) \right) \mathbf{v} n_k = 0 \end{aligned} \quad (10)$$

or in simplified form

$$\begin{aligned} \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^{n+1} \mathbf{v} + \frac{\Delta t}{2} \int_{\Omega} \mathbf{A}(\mathbf{u}^{n+1}) \operatorname{div} \mathbf{f}^{n+1} \cdot \nabla \mathbf{v} - \frac{\Delta t}{2} \int_{\partial \Omega} \mathbf{A}(\mathbf{u}^{n+1}) \operatorname{div} \mathbf{f}^{n+1} \cdot \mathbf{n} \mathbf{v} = \\ = \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^n \mathbf{v} + \int_{\Omega} \mathbf{f}^n \cdot \nabla \mathbf{v} - \int_{\partial \Omega} \mathbf{f}^n \cdot \mathbf{n} \mathbf{v} \end{aligned} \quad (11)$$

Now, using relations

$$\begin{aligned} \mathbf{f} &= \mathbf{A} \mathbf{u} \\ \operatorname{div} \mathbf{f} &= \mathbf{A} \cdot \nabla \mathbf{u} \end{aligned}$$

we can rewrite (11) in the following form

$$\begin{aligned} \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^{n+1} \mathbf{v} + \frac{\Delta t}{2} \int_{\Omega} \mathbf{A}^{n+1} (\mathbf{A}^{n+1} \cdot \nabla \mathbf{u}^{n+1}) \cdot \nabla \mathbf{v} - \frac{\Delta t}{2} \int_{\partial \Omega} \mathbf{A}^{n+1} (\mathbf{A}^{n+1} \cdot \nabla \mathbf{u}^{n+1}) \cdot \mathbf{n} \mathbf{v} = \\ = \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^n \mathbf{v} + \int_{\Omega} \mathbf{A}^n \mathbf{u}^n \cdot \nabla \mathbf{v} - \int_{\partial \Omega} \mathbf{A}^n \mathbf{u}^n \cdot \mathbf{n} \mathbf{v} \end{aligned} \quad (12)$$

Let me rewrite (12) in detail for implementation purposes

$$\begin{aligned}
& \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^{n+1} \mathbf{v} + \\
& \frac{\Delta t}{2} \int_{\Omega} A_1^{n+1} \left(A_1^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial x} + A_2^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial y} \right) \frac{\partial \mathbf{v}}{\partial x} + A_2^{n+1} \left(A_1^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial x} + A_2^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial y} \right) \frac{\partial \mathbf{v}}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\Omega} A_1^{n+1} \left(A_1^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial x} + A_2^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial y} \right) n_x \mathbf{v} + A_2^{n+1} \left(A_1^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial x} + A_2^{n+1} \frac{\partial \mathbf{u}^{n+1}}{\partial y} \right) n_y \mathbf{v} \\
& = \frac{1}{\Delta t} \int_{\Omega} \mathbf{u}^n \mathbf{v} + \int_{\Omega} \left(A_1^n \mathbf{u}^n \frac{\partial \mathbf{v}}{\partial x} + A_2^n \mathbf{u}^n \frac{\partial \mathbf{v}}{\partial y} \right) - \int_{\partial \Omega} (A_1^n \mathbf{u}^n n_x \mathbf{v} + A_2^n \mathbf{u}^n n_y \mathbf{v})
\end{aligned}$$

Therefore 1-th equation is

$$\begin{aligned}
& \frac{1}{\Delta t} \int_{\Omega} u_1^{n+1} v_1 \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_1^{n+1})_{1j} \left(\sum_k (A_1^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial x} + \sum_k (A_2^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial y} \right) \frac{\partial v_1}{\partial x} \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_2^{n+1})_{1j} \left(\sum_k (A_1^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial x} + \sum_k (A_2^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial y} \right) \frac{\partial v_1}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\partial \Omega} \sum_j (A_1^{n+1})_{1j} \left(\sum_k (A_1^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial x} + \sum_k (A_2^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial y} \right) n_x v_1 \\
& - \frac{\Delta t}{2} \int_{\partial \Omega} \sum_j (A_2^{n+1})_{1j} \left(\sum_k (A_1^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial x} + \sum_k (A_2^{n+1})_{jk} \frac{\partial u_k^{n+1}}{\partial y} \right) n_y v_1 \\
& = \frac{1}{\Delta t} \int_{\Omega} u_1^n v_1 + \int_{\Omega} \left(\sum_k (A_1^n)_{1k} u_k^n \frac{\partial v_1}{\partial x} + \sum_k (A_2^n)_{1k} u_k^n \frac{\partial v_1}{\partial y} \right) \\
& \quad - \int_{\partial \Omega} \left(\sum_k (A_1^n)_{1k} u_k^n n_x v_1 + \sum_k (A_2^n)_{1k} u_k^n n_y v_1 \right)
\end{aligned}$$

Separated into 4 parts that correspond to blocks in stiffness matrix (superscript $n + 1$ is omitted)

$$\begin{aligned}
& \frac{1}{\Delta t} \int_{\Omega} u_1 v_1 \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_1)_{1j} \left((A_1)_{j1} \frac{\partial u_1}{\partial x} + (A_2)_{j1} \frac{\partial u_1}{\partial y} \right) \frac{\partial v_1}{\partial x} + \sum_j (A_2)_{1j} \left((A_1)_{j1} \frac{\partial u_1}{\partial x} + (A_2)_{j1} \frac{\partial u_1}{\partial y} \right) \frac{\partial v_1}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\partial\Omega} \left[\sum_j (A_1)_{1j} \left((A_1)_{j1} \frac{\partial u_1}{\partial x} + (A_2)_{j1} \frac{\partial u_1}{\partial y} \right) n_x v_1 + \sum_j (A_2)_{1j} \left((A_1)_{j1} \frac{\partial u_1}{\partial x} + (A_2)_{j1} \frac{\partial u_1}{\partial y} \right) n_y v_1 \right] \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_1)_{1j} \left((A_1)_{j2} \frac{\partial u_2}{\partial x} + (A_2)_{j2} \frac{\partial u_2}{\partial y} \right) \frac{\partial v_1}{\partial x} + \sum_j (A_2)_{1j} \left((A_1)_{j2} \frac{\partial u_2}{\partial x} + (A_2)_{j2} \frac{\partial u_2}{\partial y} \right) \frac{\partial v_1}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\partial\Omega} \left[\sum_j (A_1)_{1j} \left((A_1)_{j2} \frac{\partial u_2}{\partial x} + (A_2)_{j2} \frac{\partial u_2}{\partial y} \right) n_x v_1 + \sum_j (A_2)_{1j} \left((A_1)_{j2} \frac{\partial u_2}{\partial x} + (A_2)_{j2} \frac{\partial u_2}{\partial y} \right) n_y v_1 \right] \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_1)_{1j} \left((A_1)_{j3} \frac{\partial u_3}{\partial x} + (A_2)_{j3} \frac{\partial u_3}{\partial y} \right) \frac{\partial v_1}{\partial x} + \sum_j (A_2)_{1j} \left((A_1)_{j3} \frac{\partial u_3}{\partial x} + (A_2)_{j3} \frac{\partial u_3}{\partial y} \right) \frac{\partial v_1}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\partial\Omega} \left[\sum_j (A_1)_{1j} \left((A_1)_{j3} \frac{\partial u_3}{\partial x} + (A_2)_{j3} \frac{\partial u_3}{\partial y} \right) n_x v_1 + \sum_j (A_2)_{1j} \left((A_1)_{j3} \frac{\partial u_3}{\partial x} + (A_2)_{j3} \frac{\partial u_3}{\partial y} \right) n_y v_1 \right] \\
& + \frac{\Delta t}{2} \int_{\Omega} \sum_j (A_1)_{1j} \left((A_1)_{j4} \frac{\partial u_4}{\partial x} + (A_2)_{j4} \frac{\partial u_4}{\partial y} \right) \frac{\partial v_1}{\partial x} + \sum_j (A_2)_{1j} \left((A_1)_{j4} \frac{\partial u_4}{\partial x} + (A_2)_{j4} \frac{\partial u_4}{\partial y} \right) \frac{\partial v_1}{\partial y} \\
& - \frac{\Delta t}{2} \int_{\partial\Omega} \left[\sum_j (A_1)_{1j} \left((A_1)_{j4} \frac{\partial u_4}{\partial x} + (A_2)_{j4} \frac{\partial u_4}{\partial y} \right) n_x v_1 + \sum_j (A_2)_{1j} \left((A_1)_{j4} \frac{\partial u_4}{\partial x} + (A_2)_{j4} \frac{\partial u_4}{\partial y} \right) n_y v_1 \right] \\
& = \frac{1}{\Delta t} \int_{\Omega} u_1^n v_1 + \int_{\Omega} \left(\sum_k (A_1^n)_{1k} u_k^n \frac{\partial v_1}{\partial x} + \sum_k (A_2^n)_{1k} u_k^n \frac{\partial v_1}{\partial y} \right) \\
& \quad - \int_{\partial\Omega} \left(\sum_k (A_1^n)_{1k} u_k^n n_x v_1 + \sum_k (A_2^n)_{1k} u_k^n n_y v_1 \right)
\end{aligned}$$